#### Vehicle Dynamics and Simulation

## Using Eigenvalues and Eigenvectors

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#### Lecture overview

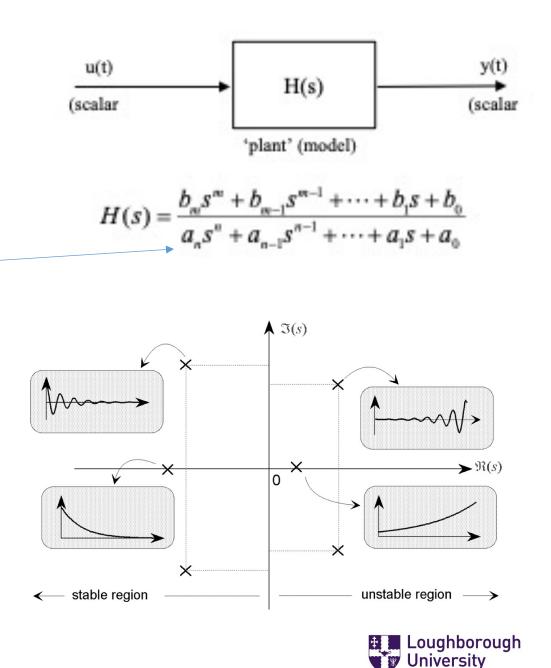
- Transfer functions
- Modal motion in free vibration
  - Eigenvalues
  - Eigenvectors





## **Transfer Functions**

- Transfer functions relate input to output
- The roots/poles of the characteristic equation determine <u>frequency and</u> <u>damping of each mode</u> i.e. the dynamics of the system
- In state space form additional information is also available describing <u>mode shapes</u> from the A matrix

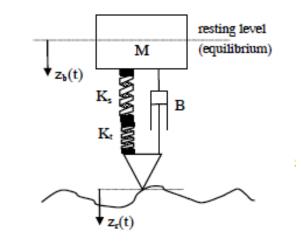


## Simple Example

- Using the simple suspension example from Section 2b
  - Assume zero initial conditions
  - Take Laplace transform
  - Write transfer function
  - Enter parameter values
- The roots of the characteristic equation in this example are;

$$-1.875 \pm 6.372i$$

• The nature of the roots e.g. complex, repeated, distinct and real determine the general solution approach. They also define the dynamics of the system.



$$H(s) = \frac{Y(s)}{U(s)} = \frac{Bs + K}{Ms^2 + Bs + K}$$

 $H(s) = \frac{3.75s + 44.1}{s^2 + 3.75s + 44.1}$ 



## Laplace Transform and the Transfer Function

• State space representation;

$$\dot{\boldsymbol{x}} = A\boldsymbol{x} + B\boldsymbol{u}$$
 [1]

$$y = Cx + Du$$
[2]

• Assuming zero initial conditions and taking the Laplace transform of [1];

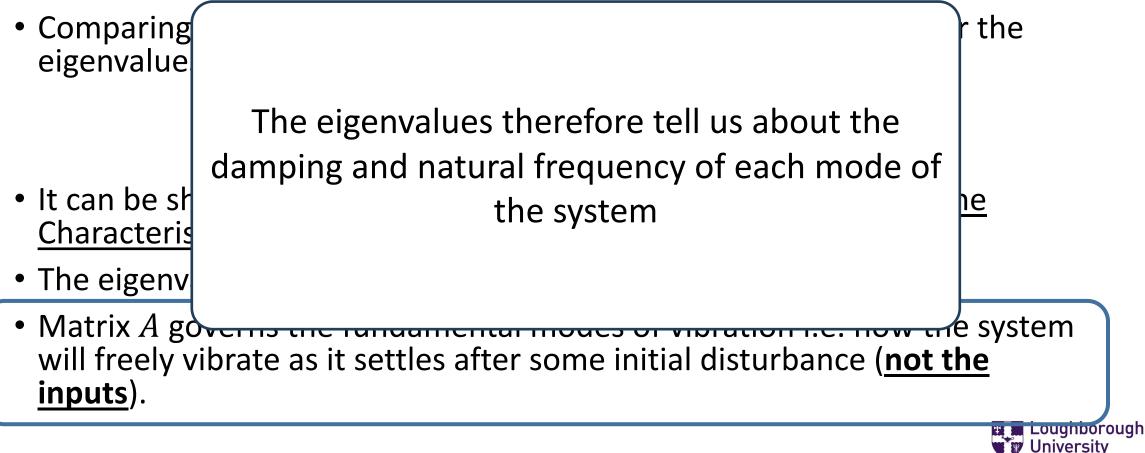
sX = AX + BU(sI - A)X = BU $X = (sI - A)^{-1}BU$  $Y = C(sI - A)^{-1}BU + DU$  $H(s) = C(sI - A)^{-1}B + D$ 

[3]

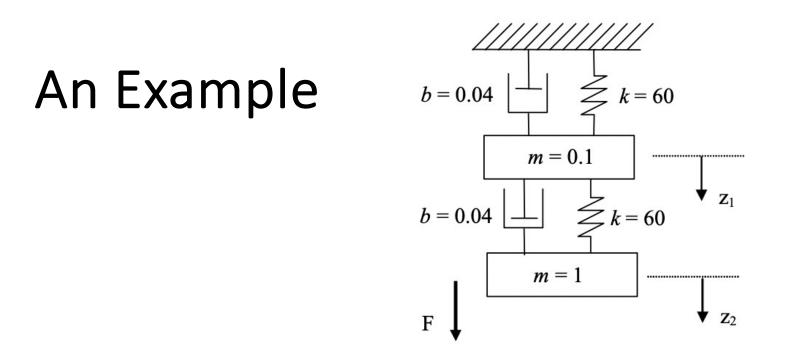
• Substituting into [2];

# Laplace Transform and the Transfer Function

• Equation [3] provides a general solution in terms of the transfer function, H(s) and is an <u>alternate form</u> to the State Space Representation.



## worked example (reminder)



• The equations of the system are (expressed in terms of two second order differential equations);

$$M_2 \ddot{z}_2 = F - k(z_2 - z_1) - b(\dot{z}_2 - \dot{z}_1)$$
$$M_1 \ddot{z}_1 = k(z_2 - z_1) + b(\dot{z}_2 - \dot{z}_1) - kz_1 - b\dot{z}_1$$



Choosing one deflection and one velocity state per mass;

and the input, u = F. So that the set of equations describing the state derivatives becomes;

 $\dot{x}_1 = x_3$  $\dot{x}_2 = x_4$  $\dot{x}_3 = -\frac{2k}{m}x_1 + \frac{k}{m}x_2 - \frac{2b}{m}x_3 + \frac{b}{m}x_4$  $\dot{x}_4 = kx_1 - kx_2 + bx_3 - bx_4 + u$ 

So that the state space representation becomes;

$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \\ \dot{x}_3 \\ \dot{x}_4 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -2k/m & k/m & -2b/m & b/m \\ k & -k & b & -b \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} u$$
$$\dot{X} = AX + BU$$

[3]

#### Modal motion in free vibration – Eigenvalues

• The vector of deflections only,  $z(t) = [x_1 \ x_2]^T$  for our example may be written as a linear combination;

$$\boldsymbol{z}(\boldsymbol{t}) = Re\{\boldsymbol{u}_1 e^{\lambda_1 t} + \boldsymbol{u}_2 e^{\lambda_2 t}\}$$

where each term  $u_i e^{\lambda_i t}$  represents a single vibrational mode,  $u_i$  are complex constants [2x1 vector in this example],  $\lambda_i$  are complex scalars and  $\max(i) = n$  with n being the number of states.

• Evaluating a single term in the above, split  $\lambda_i$  into real and imaginary parts;

$$\lambda_i = \sigma + bi$$

• Using the above and Euler's formula we can better evaluate what is happening;

$$\boldsymbol{u}_i e^{\lambda_i t} = \boldsymbol{u}_i e^{(\sigma + bi)t} = \boldsymbol{u}_i e^{\sigma t} e^{ibt}$$

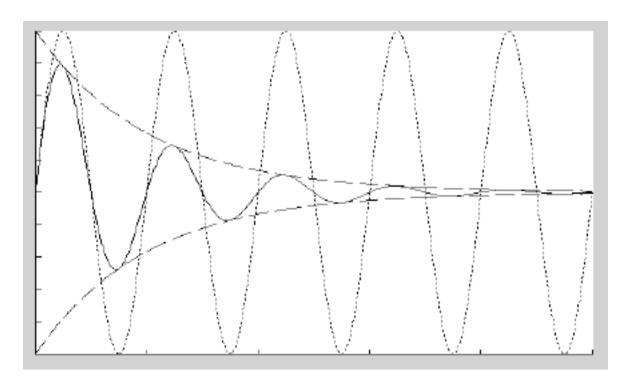
 $\boldsymbol{u}_i e^{\lambda_i t} = \boldsymbol{u}_i e^{\sigma t} (\cos(\boldsymbol{b}t) + i\sin(\boldsymbol{b}t))$ 

[5]



#### Modal motion in free vibration - Eigenvalues

• From [5],  $\sigma$  should be *-ve* bounding the response to a decaying exponential, b gives the frequency of the sinusoidal component,  $u_i$  (complex) determines the magnitude and the relative phase of each mode.



$$\boldsymbol{u}_i e^{\lambda_i t} = \boldsymbol{u}_i e^{\sigma t} (\cos(\boldsymbol{b}t) + i\sin(\boldsymbol{b}t))$$

Modal decomposition of response Solid line = total response Short dash = sinusoidal component Long dash = exponential decay



#### Modal motion in free vibration – Eigenvalues

• Eigenvalues appear in (complex conjugate) pairs and can be written;

$$\lambda_{1,2} = \sigma \pm j \omega_d$$

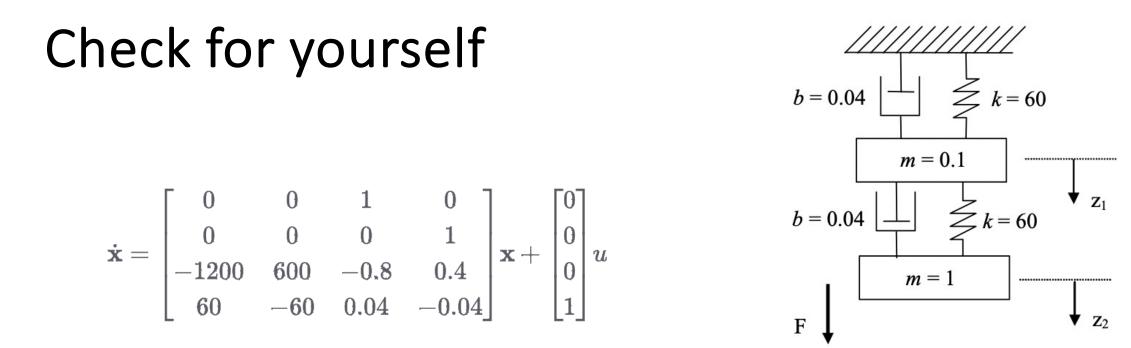
where  $\sigma$  is the modal damping factor and  $\omega_d$  is the damped natural frequency.



## Check for yourself

- Use the eig() function in MATLAB to determine the eigenvalues of matrix A from the previous example.
- How are the complex conjugate pairs placed within the resulting vector?





 $\lambda_{1,2} = \sigma \pm j\omega_d$ 

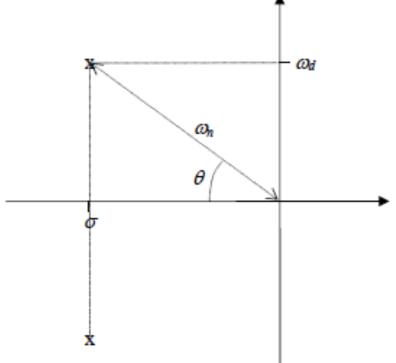
D =

-0.41025 +	35.08i	0 +	Øi	0 +	0i	0 +	0i
0 +	0i	-0.41025 -	35 <b>.</b> 08i	0 +	0i	0 +	0i
0 +	0i	0 +	0i	-0.0097502 +	5.4084i	0 +	0i
0 +	0i	0 +	Øi	0 +	Øi	-0.0097502 -	5 <b>.</b> 4084i



## Modal motion in free vibration – Eigenvalues

- From the eigenvalues it is possible to tell
  - Damped natural frequency,  $\omega_d$
  - Natural frequency [Hz],  $\omega_n/2\pi$
  - Damping factor,  $\sigma$
  - Damping ratio,  $\zeta = \cos(\theta)$
  - Settling time (within 2%),  $T_s = \frac{4}{|\sigma|}$
  - Percent overshoot,  $100e^{\frac{-n\zeta}{\sqrt{1-\zeta^2}}}$
- Note:  $\lambda = 0$  corresponds to the steady-state response of the system (not dynamics)





## Check for yourself

- Using the previous example find the eigenvalues of the system and hence determine;
  - Damped natural frequency,  $\omega_d$
  - Natural frequency [Hz],  $\omega_n/2\pi$
  - Damping factor,  $\sigma$
  - Damping ratio,  $\zeta = \cos(\theta)$
  - Settling time (within 2%),  $T_s = \frac{4}{|\sigma|}$
  - Percent overshoot,  $100e^{\frac{-\pi\zeta}{\sqrt{1-\zeta^2}}}$



#### Modal motion in free vibration – Eigenvectors

- Eigenvectors can show the magnitudes at which the states vibrate in relation to one another.
- Writing eigenvalues and eigenvectors together in matrix form;

AV = VD

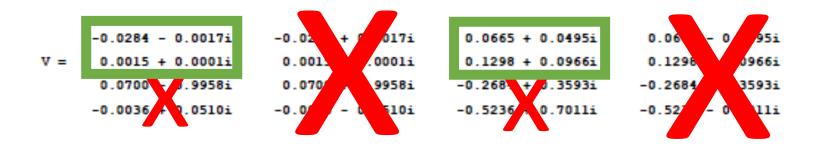
where;

$$\boldsymbol{V} = \begin{bmatrix} u_1 & u_2 & \cdots & u_n \\ \lambda_1 u_1 & \lambda_2 u_2 & \cdots & \lambda_n u_n \end{bmatrix} \text{ and } \boldsymbol{D} = \begin{bmatrix} \lambda_1 & & & \\ & \lambda_2 & & \\ & & \ddots & \\ & & & \ddots & \\ & & & & & \lambda_n \end{bmatrix}$$



## Modal motion in free vibration – Eigenvectors

 Using MATLAB 'eig(A)' to find the eigenvectors of the example system, A matrix;

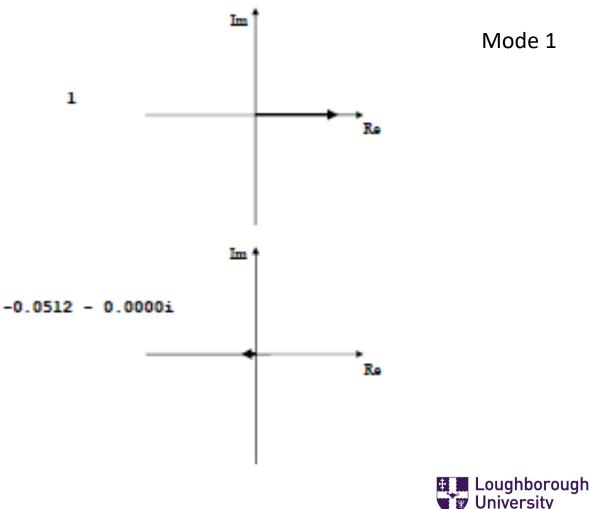


- Note
  - The second and fourth columns are the complex conjugates of the first and third columns respectively
  - Rows three and four are the first and second rows multiplied by their respective eigenvalues
  - The system can then be characterized by considerably less 'unique information'



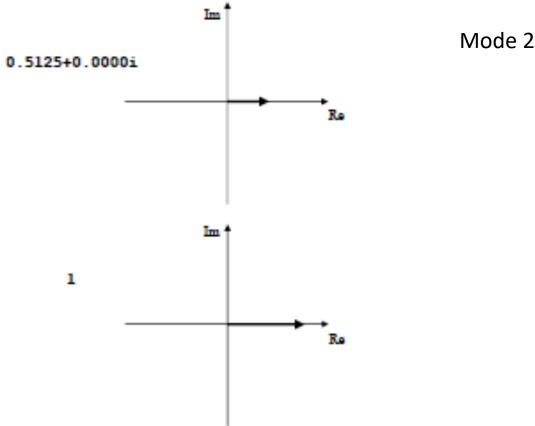
## Modal motion in free vibration – Eigenvectors

- Dividing through by the largest magnitude eigenvector (-0.0284 0.017*i*) to normalize the eigenvectors.
- Plot the eigenvector components (first mode)
- The relative magnitude and phase is seen on the two plots



## Modal motion in free vibration – Eigenvectors

- Similarly for the second (non-conjugate) mode of interest
- The relative magnitude and phase is seen on the two plots
- Note the differences between first and second modes of vibration





## Conclusions

- Transfer function vs state space representation
- Eigenvalues tell us;
  - Damped natural frequency
  - Natural frequency
  - Damping factor
  - Damping ratio
  - Settling time
  - Percent overshoot
- Eigenvectors help us to understand vibration of the modes relative to one another

